APPROXIMATE MAXIMUM LIKELIHOOD ESTIMATION FOR A SPATIAL POINT PATTERN[*]

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Several authors have proposed stochastic and non-stochastic approximations to the maximum likelihood estimate for a spatial point pattern. This approximation is necessary because of the difficulty of evaluating the normalizing constant. However, it appears to be neither a general theory which provides grounds for preferring a particular method, nor any extensive empirical comparisons. In this paper, we review five general methods based on approximations to the maximum likelihood estimate which have been proposed in the literature. We also present the results of a comparative simulation study developed for the Strauss model.

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1. INTRODUCTION

A spatial point pattern is a set of points

$$X = \{x_i \in A : i = 1, \ldots, n\}$$

for some planar region A. The x_i are called events to distinguish them from generic points $x \in A$. Very often, A is a sampling window within a much larger region and it is reasonable to regard X as a partial realization of a planar point process, the events consisting of all points of the process which lie within A.

Parameter estimation for two-dimensional point pattern data is difficult, because most of the available stochastic models have intractable likelihoods (see Ripley, 1977, 1988 and Diggle, 1983). An exception is the class of Gibbs or Markov point processes (Baddeley and Moller, 1989; Ripley, 1989), where the likelihood $l(X; \theta)$ typically forms an exponential family and is given explicitly up to a normalizing constant. However, the latter is not known analytically precluding the use of exact maximum likelihood, so parameter estimates must be based on approximations.

Gibbs point processes first appeared in the theory of statistical physics, where Gibbs distributions were applied to describe the equilibrium states of closed physical systems of interacting objects. In mathematical statistics Gibbs point processes are used as models of spatial point patterns. A preliminary paper introducing the Gibbs processing into the statistical literature is Ripley and Kelly (1977). Examples can be found in biology, plant ecology, forestry and economy.

The topic of this paper concerning Gibbs type processes has a general validity arising from two aspects: (i) It is a general way of proceeding in cases of exponential families with dependent samples, and (ii) it has theoretical value on its own. Examples of (i) are the applications of Markov random fields for lattice data (Besag, 1974; Geyer and Thompson, 1992), Markov random fields in image analysis (Geman and Geman, 1984), Gibbs point processes and germ-grain models in high level image analysis (Baddeley and van Lieshout, 1993), modelling of random graphs and general interaction models (Strauss, 1986). Gibbs processes are useful as prior distributions in image interpretation tasks, such as object recognition, edge detection and feature extraction (van Lieshout and Baddeley, 1995; Molina and Ripley, 1989). Maximum likelihood solutions tend to suffer from multiple response and the prior distribution serves to penalize scenes with too many almost identical objects, disconnected or crossing edges. Usually, the posterior distribution also possesses a Markov property, enabling sampling and optimization by iterative procedures that recursively update the scene by simple operations of addition or deletion.

In this paper, we consider generally applicable methods for estimating the parameter θ confining our attention to stochastic and non-stochastic approximations to the ma-

ximum likelihood estimate (MLE). We use a simple point process model, the Strauss process (Strauss, 1975), to illustrate and compare these methods which could be applied to more general and complex models. The Strauss process is a point process model which has been used in modelling (non-clustered) point patterns in some of the mentioned references and is a demanding member of the exponential family for a dependent sample.

The interest of the present paper relies on methods of estimation which can be used routinely in applications, and which do not place artificial restrictions on the parametric form of $l(X; \theta)$. The aim is to present a comparative study among the approximations to the MLE and to discuss the practical implications. We consider only homogeneous, i.e., stationary and isotropic processes. Throughout this paper, $N(A)$ stands for the number of events in A, $|A|$ denotes the area of A and $\lambda = E[N(A)]/|A|$ denotes the intensity of the process.

For a general introduction to statistical methodology for spatial point patterns, see for example Ripley (1981), Diggle (1983), Stoyan, Kendall and Mecke (1995) and Cressie (1993). Other parametric methods of estimation, not considered here, are maximum pseudo-likelihood and the Takacs-Fiksel method (Diggle et al., 1994; Takacs, 1986). In a different vein, Diggle, Gates and Stibbard (1987) develop a smooth, non-parametric estimator for the interaction function, to which a parametric family could be fitted by standard curve-fitting techniques such as non-linear least squares.

The plan of the paper is as follows. Section 2 describes the approximate MLE methods for a particular Gibbs process, the Strauss model. Section 3 shows the simulation study to compare the different methods. The paper ends with a section of final conclusions.

2. APPROXIMATE MLE FOR A GIBBS PROCESS

A class of stochastic models for patterns of n events in a bounded region A is the class of *pairwise interaction point processes*. The joint density for a pattern X, taken with respect to the Poisson measure μ, is given by

$$(1) \qquad f(X; \theta) = C(\theta)^{-1} \beta^n \exp \left\{ -\sum_{i=1}^{n} \sum_{j>i} \Phi(\|x_i - x_j\|; \theta) \right\} / n!$$

In (1) , $\|.\|$ denotes Euclidean distance, $\Phi(.)$ is a *potential function* depending on a set of parameters θ, β is a parameter which determines the intensity of the process, and $C(\theta)$ is a normalizing constant. We call $U_n(X; \theta) = \sum_{i=1}^{n} \sum_{j>i} \Phi(\|x_i - x_j\|; \theta)$ the *total potential energy*. Often, (1) is written in terms of an *interaction function* $e(t) = \exp(-\Phi(t))$. Such class of point processes belongs to a more general kind of processes called *Gibbs processes* (Kelly and Ripley, 1976; Daley and Vere-Jones, 1988; Baddeley and Moller,

1989). Note that restrictions on the form of the potential $\Phi(.)$ are needed to ensure that the normalizing constant in (1) is finite.

A *Strauss process* (Strauss, 1975) is a pairwise interaction process in which the density depends only on the number of neighbour pairs defined by

$$s(X) = \sum_{i=1}^{n} \sum_{j>i} I(\|x_i - x_j\| \leq r).$$

Considering in (1) the Strauss potential function

$$\Phi(t) = \begin{cases} -\log(\theta), t \leq r \\ 0, t > r \end{cases}$$

the likelihood takes the form (Kelly and Ripley, 1976)

$$l(X;\theta) = \exp(-|A|)\alpha(\theta)^{-1}\beta^n\theta^{s(X)}$$

where the normalizing constant is $C(\theta) = \alpha(\theta)/\exp(-|A|)n!$ The case $\theta = 1$ corresponds to a Poisson process with intensity β. If $\theta = 0$, the result is a simple inhibition process that contains no events at a distance less than or equal to r. Values of $\theta < 1$ correspond to regularity of events, whilst for $\theta > 1$ the process should result in clustering (see Figures 1a, 1b and 1c). For a clustered pattern, as was pointed out by Kelly and Ripley (1976), the condition $\theta > 1$ violates the requirement of a finite normalizing constant $C(\theta)$ in (1). This problem can be removed by conditioning to the number of events, say $N = n$. This is not an artificial restriction because $n(X)$ usually provides little information about the interactions among the events. The effect on conditioning to the MLE for the Strauss family has been demonstrated by Geyer and Moller (1994). Furthermore, conditioning on n makes it easier to generate simulations by the discrete-time Markov chain method of Ripley (1979, 1987). The conditional likelihood function for the Strauss process is given by

$$(2) \qquad\qquad l_n(X;\theta) = \theta^{s(X)}/C_n(\theta)$$

where the normalizing constant is given by

$$(3) \qquad\qquad C_n(\theta) = \int_{A^n} \theta^{s(X)} dx_1 \cdots dx_n.$$

Maximum likelihood estimation of θ requires the evaluation of (3) which is not usually obtainable in closed form. We therefore try to maximize an approximation to the likelihood function. In the following, we develop approximations to the MLE for the Strauss conditional model.

Figure 1a. Realizations of simulated patterns under the Strauss model for different values of parameter θ. In each pattern it is also included $s(X)$, the number of neighbour pairs. r=0.10

Figure 1b. Realizations of simulated patterns under the Strauss model for different values of parameter θ. In each pattern it is also included $s(X)$, the number of neighbour pairs. r=0.15
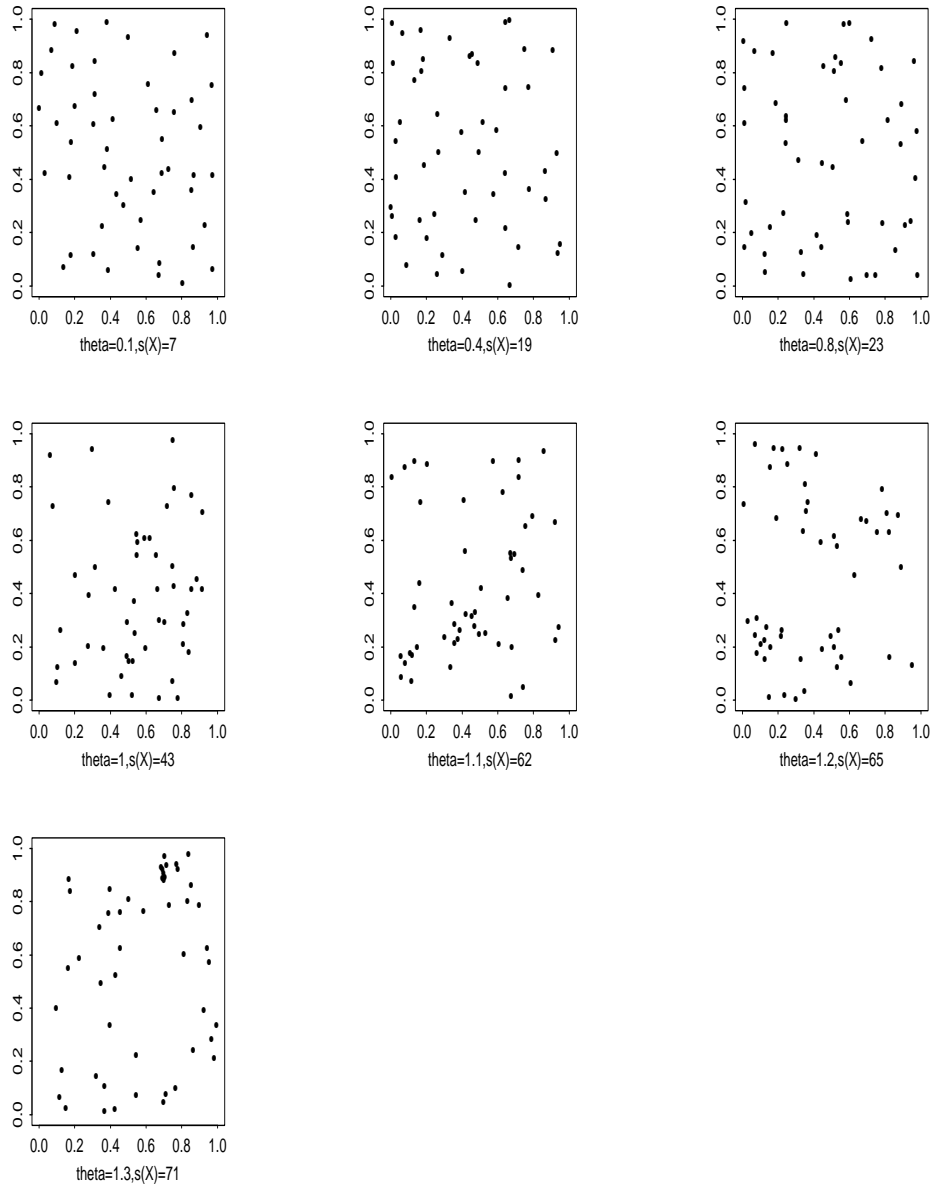
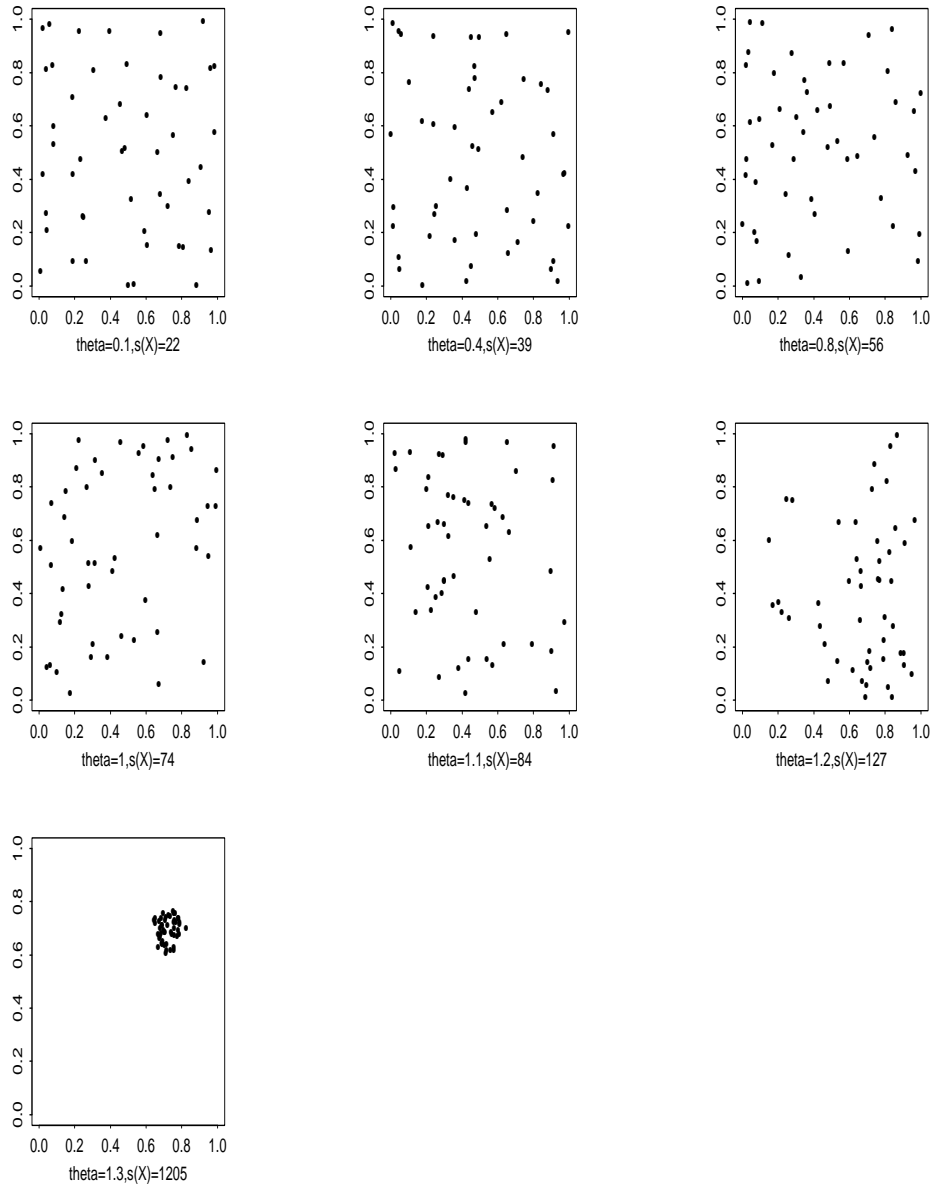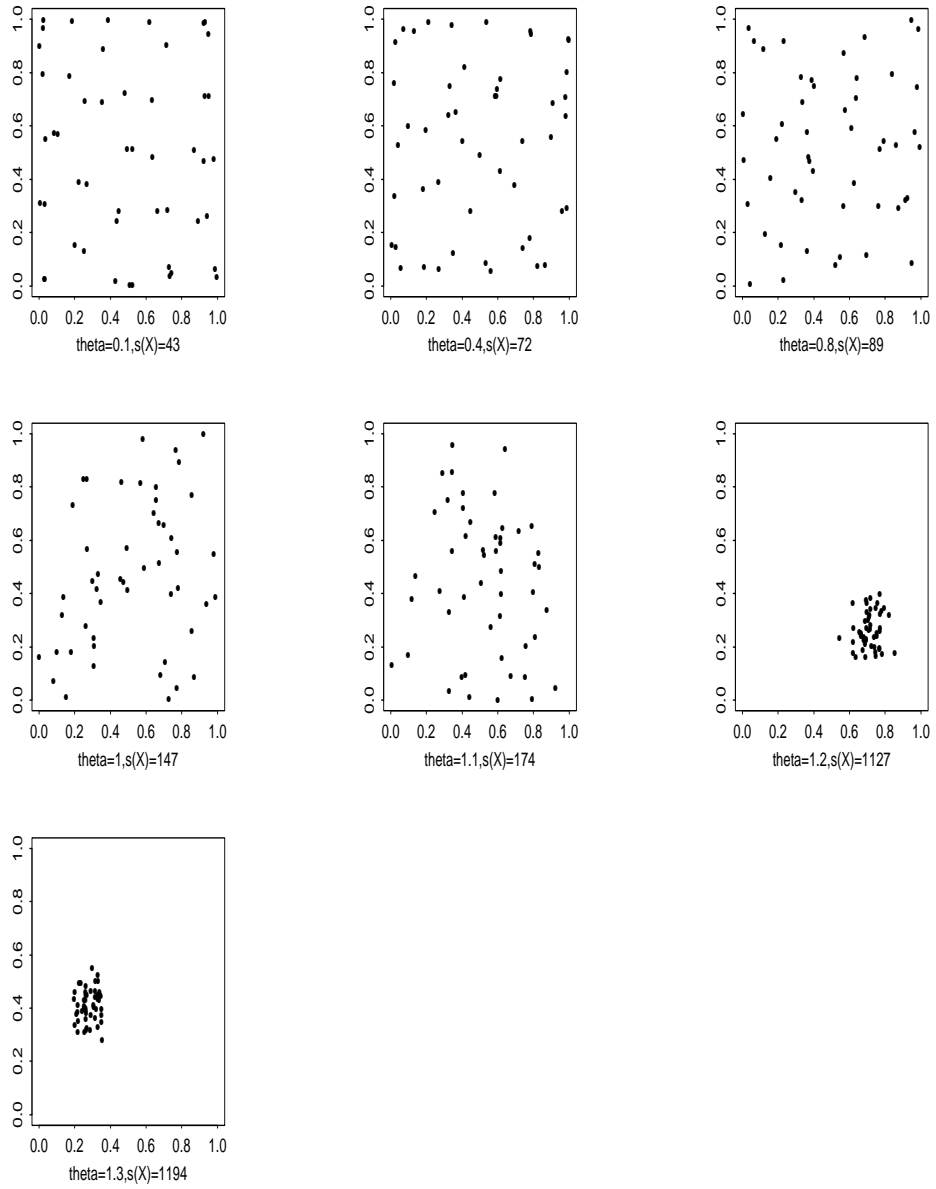**Figure 1c.** Realizations of simulated patterns under the Strauss model for different values of parameter θ. In each pattern it is also included $s(X)$, the number of neighbour pairs. r=0.20

2.1. Method of Ogata-Tanemura

Ogata and Tanemura (1981) proposed to use a cluster-expansion method of statistical mechanics assuming that the events of the point process are sparsely distributed, so that third and higher-order cluster integrals are negligible. Then, using up to the second-order cluster integral, an approximation to the normalizing constant is given by

$$C_n(\theta) = |A|^n \left\{ 1 - \frac{b(\theta)}{|A|} \right\}^{n(n-1)/2}$$

where $b(\theta)$ is, for the Strauss model, $b(\theta) = \pi(1-\theta)r^2$. Then, the MLE is given by

(4)
$$\widehat{\theta} = \frac{s(X)\left\{2|A| - \pi r^2\right\}}{\pi r^2 (n(n-1)/2 - s(X))}.$$

2.2. Method of Penttinen

Penttinen (1984) proposed another sparse-data approximation to (3), which for the Strauss model takes the form

$$C_n(\theta) = \exp\left\{1/2n(n-1)\pi r^2(\theta-1)\right\}$$

and the MLE is given by

(5)
$$\widehat{\theta} = \frac{s(X)}{1/2n(n-1)\pi r^2}.$$

2.3. Method of virial expansions

This method consists of the following approximation of (3),

(6)
$$n^{-1}\log(C_n) \approx (b_n/2)\int_{\Re^2} f_{12}dx_2 + (b_n^2/4)\int_{\Re^4} f_{12}f_{13}f_{23}dx_2dx_3 + (b_n^3/8)\int_{\Re^6}(f_{12}f_{13}f_{14}f_{23}f_{24}f_{34} + 6f_{12}f_{13}f_{14}f_{23}f_{24} + 3f_{12}f_{14}f_{23}f_{34})dx_2dx_3dx_4 + \cdots$$

where $b_n = n/|A|$ and $f_{ij} = \exp(-\Phi(\|x_i - x_j\|;\theta)) - 1$ (Ripley, 1988). To implement this method for the Strauss process we use the fourth order expansion obtained by calculating the integrals in (6):

$$\log(C_n(\theta)) = -\pi n(n-1)\Psi r^2/(2|A|) - 0{,}29325\pi^2 \frac{n!}{(n-3)!}\Psi^3 r^4/(6|A|^2)$$
$$-\pi^3 \frac{n!}{(n-4)!}\left\{-0{,}27432\Psi^6 + 2{,}18542\Psi^5 - 1{,}37886\Psi^4\right\}r^6/(24|A|^3)$$

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where $\Psi = 1 - \theta$.

Then solving

(7)
$$d\log(C_n(\theta))/d\theta = s(X)/\theta$$

we obtain $\widehat{\theta}$, the approximate MLE of θ.

2.4. Stochastic approximation based on a Newton-Raphson procedure

Penttinen (1984) suggested a Newton-Raphson type algorithm for solving the maximum likelihood estimating equation. Assume $\Phi(t;\theta)$ is twice differentiable with respect to θ. Differentiation of both sides of equation (3) yields

$$\frac{-\partial C_n(\theta)}{\partial \theta} = C_n(\theta)E_\theta[\partial U_n(x_1,\dots,x_n;\theta)/\partial \theta]$$

where the *total potential energy*, for the Strauss process, is

$$U_n(x_1,\dots,x_n;\theta) = -\log\theta^{s(X)}.$$

The MLE $\widehat{\theta}$ solves $\partial \log(l_n(X;\theta))/\partial\theta = 0$. If $\widehat{\theta}_0$ denotes an initial guess for $\widehat{\theta}$, then the Newton-Raphson algorithm consists of

$$\widehat{\theta}_{k+1} = \widehat{\theta}_k - \left[\overline{\Gamma}_T(\widehat{\theta}_k)\right]^{-1}\overline{\beta}_T(\widehat{\theta}_k) \quad k = 0,1,2,\dots$$

where

(8)
$$\overline{\beta}_T(\widehat{\theta}_k) = \frac{1}{T}\sum_{t=1}^{T}\frac{1}{\widehat{\theta}_k}[s(X) - s(\phi_n(t))]$$

and

$$\overline{\Gamma}_T(\widehat{\theta}_k) = \frac{1}{T}\sum_{t=1}^{T}\frac{1}{\widehat{\theta}_k^2}[s(\phi_n(t)) - s(X)]$$
$$- \left\{\frac{1}{\widehat{\theta}_k}[s(X) - s(\phi_n(t))] - \overline{\beta}_T(\widehat{\theta}_k)\right\}^2.$$

Note that $\phi_n(1),\dots,\phi_n(T)$ are simulated according to a Strauss process with parameter $\widehat{\theta}_k$.

2.5. Stochastic approximation based on Robbins-Monro procedure

This stochastic approximation procedure was first introduced by Robbins and Monro (1951) and can be used to estimate the solution θ^* of an equation $F(\theta^*) = \varphi$ when

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there is very little information about the function F but it is possible, for any given θ, to generate a random variable T_θ with expectation $E(T_\theta) = F(\theta)$.

For the Strauss model, the goal is to solve

(9)
$$M(\widehat{\theta}) = s(X)$$

for $\widehat{\theta}$, where X is the observed data and $M(\theta) = E_\theta[s(X)]$. Then we set $T_\theta = s(X_\theta)$, where X_θ is a simulated Strauss process with parameter θ and we obtain, recursively, a sequence of estimates of $\widehat{\theta}$ using

$$\theta_{k+1} = \theta_k + \frac{B}{k}\left\{s(X) - s(X_{\theta_k})\right\}.$$

Then $\theta_k \to \widehat{\theta}$ (a.s.) (Moyeed and Baddeley, 1991).

Defining $\mu = M'(\widehat{\theta})$ and $\sigma^2 = Var_\theta[s(X)]$, if $B > 1/(2\mu)$ then θ_k is asymptotically normally distributed with mean $\widehat{\theta}$ and variance $B^2\sigma^2/(2B\mu - 1)$.

The starting value θ_0 is arbitrary, but should be set to an initial approximation such as that holding in the sparse case

$$\theta_0 = \frac{2s(X)|A|}{n(n-1)\pi r^2}.$$

The optimum B, B_{opt}, could be estimated by

$$B_{opt} = \frac{1}{\mu} = \frac{1}{M'(\widehat{\theta})}$$

or

$$B_{opt} = \frac{2|A|}{n(n-1)\pi r^2}.$$

3. A SIMULATION STUDY

3.1. Edge-correction

Commonly, the region A is a sampled sub-region of a much larger region within which the phenomenon operates and some form of edge-correction is vital. When A is a rectangle, a possible strategy is to map A onto a torus by identifying opposite edges. This periodic boundary is commonly used for computer experiments in statistical mechanics.

However, for the analysis of real data, periodic boundaries can introduce undesirable artefacts: toroidal distances can be arbitrarily small even when the underlying process has a positive hard-core distance. In the present comparative simulation study, the points patterns were themselves generated using a periodic boundary, then this particular difficulty does not arise.

To compensate for the omission of contributions to the total potential from unobserved events outside A we replace summations of the form

$$\sum_{j>i} \Phi(\|x_i - x_j\|; \theta)$$

by

$$\frac{1}{2} \sum_{j \neq i} w_{ij}^{-1} \Phi(\|x_i - x_j\|; \theta)$$

where w_{ij} is the proportion of the circumference of the circle with centre x_i and radius $\|x_i - x_j\|$ which is contained within A. This is an adapted version of Ripley's correction (Ripley, 1977, 1988). The majority of available edge-corrections correct the bias using lengths or areas of parts of circles or discs, respectively.

In the simulation study, we also include results using the so-called free boundary conditions, in which no edge-correction at all is made.

3.2. Standard Errors

One possible way to obtain approximate standard errors is by using Monte Carlo methodology. For this approach, we simulate s realisations with $\theta = \widehat{\theta}$, the point estimate under the chosen method for the original data. We then evaluate point estimates $\widehat{\theta}_j, j = 1, \ldots, s$ from the simulated patterns and use the empirical distribution of the $\widehat{\theta}_j$ as an approximation to the sampling distribution of $\widehat{\theta}$. In particular, the sample mean and standard deviation of the $\widehat{\theta}_j$ give useful indications of the bias and efficiency of estimation. This Monte Carlo approach is highly computer-intensive and it is usually known as parametric bootstrap.

3.3. Simulation method

The *spatial birth-and-death process* provides the framework under which Ripley (1977, 1979) proposes to simulate a Markov point process on the bounded Borel set $A \subset \mathfrak{R}^d$ with n fixed. The method is related to Markov processes used in statistical mechanics and surveyed by Hastings (1970). Consider a set of particles interacting according to a certain potential function on a set A with periodic boundary, i.e. A is identified

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with a torus. First, select n events from a uniform distribution on A; call this initial point pattern $\phi_n(0)$. At step $(t+1)$, delete systematically in turn one of the n events of $\phi_n(t) = \{x_1, \ldots, x_n\}$, say event x_i, and let $\phi_n(t) - \{x_i\}$ denote the point pattern formed by removing x_i from $\phi_n(t)$. Let

$$p(u; \phi_n(t) - \{x_i\}) = \frac{l_n(\phi_n(t) - \{x_i\}, u)}{l_{n-1}(\phi_n(t) - \{x_i\})}$$

denote the conditional intensity at $u \in A$ given $\phi_n(t) - \{x_i\}$. Define

$$M = \sup_{u \in A} p(u; \phi_n(t) - \{x_i\}).$$

Select an event u from a uniform distribution on A and set $\phi_n(t+1) = \{\phi_n(t) - \{x_i\}, u\}$ with probability $p(u; \phi_n(t) - \{x_i\})/M$; otherwise, selection is repeated until a qualifying u is found. This method ensures that samples taken every n steps have no points in common. Ultimately, convergence to a Markov point process with likelihood $l_n(.)$ will occur.

Unfortunately, in the case of the Strauss model, for θ much larger than 1 the algorithm is very slow and may result in simulation difficulties (see Figures 1b and 1c when $\theta = 1,2$ and 1,3).

3.4. Design of the Simulation Study

For the simulation study we selected eight parameter values: $\theta = 0,1,0,4$ and $0,8$ corresponding to regular patterns; $\theta = 1$ for the random pattern (Poisson process) and $\theta = 1,1,1,2$ and $1,3$ for clustered ones (strongly interactive patterns). We also considered three different ranges of interaction: $r = 0,1,0,15$ and $0,2$. For each combination of parameter value and range of interaction we simulated 100 realizations, each one with $n = 50$ events on A the unit square. From the simulated realization we evaluated the estimate of θ using the five methods of approximation described in Section 2 and incorporating the edge-correction described in Section 3.1.

Each combination of parameter value, interaction range, method of estimation and edge-correction (no edge-correction, Ripley's and toroidal (periodic)) therefore yielded 100 estimates $\widehat{\theta}_j, j = 1, \ldots, 100$, which are summarised by the box-plots shown in Figures 2a, 2b and 2c.

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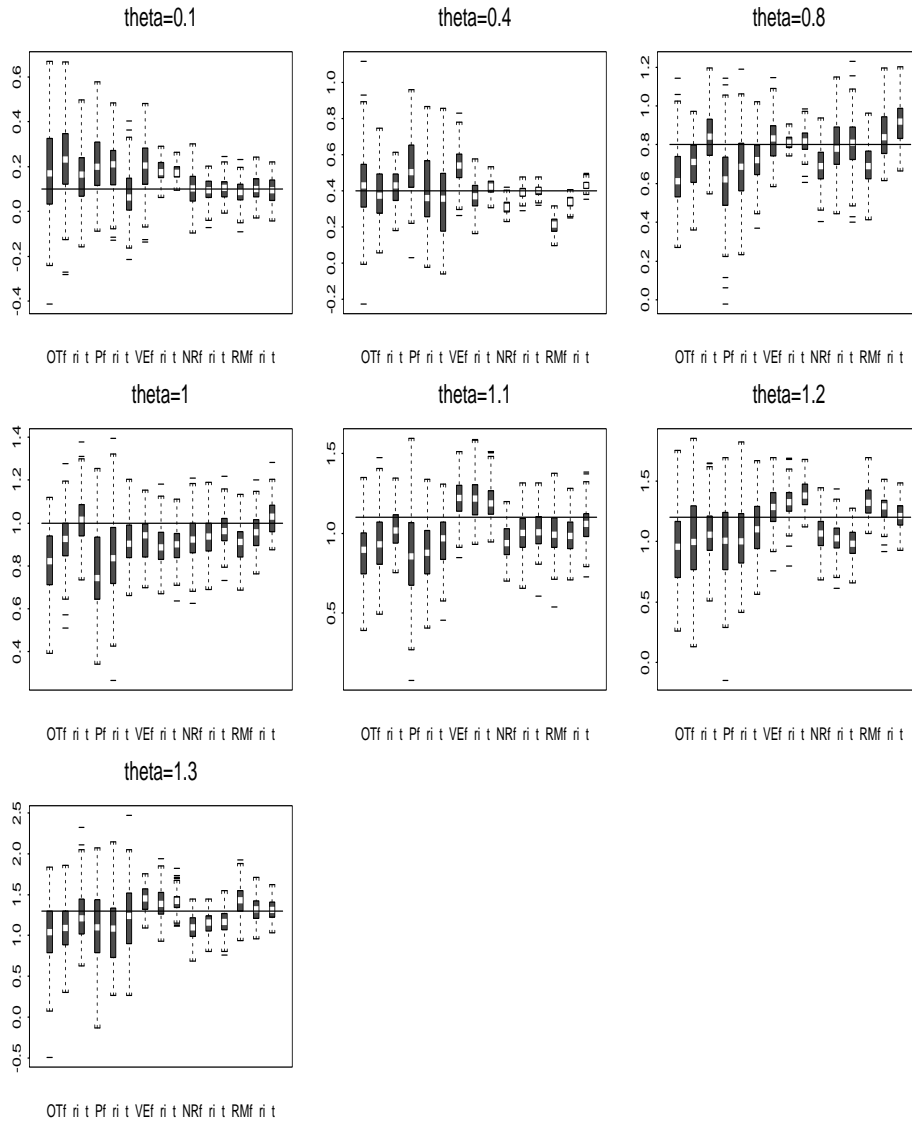
Figure 2a. Box-plots of simulated parameter estimates. The horizontal lines indicate the true value of θ. The upper case letter identifies the method of estimation (OT=Ogata-Tanemura, P=Penttinen, VE=Virial Expansions, NR=Newton-Raphson, RM=Robbins-Monro), the lower case letter identifies the boundary condition (f=free, ri=Ripley, t=toroidal). r=0.10

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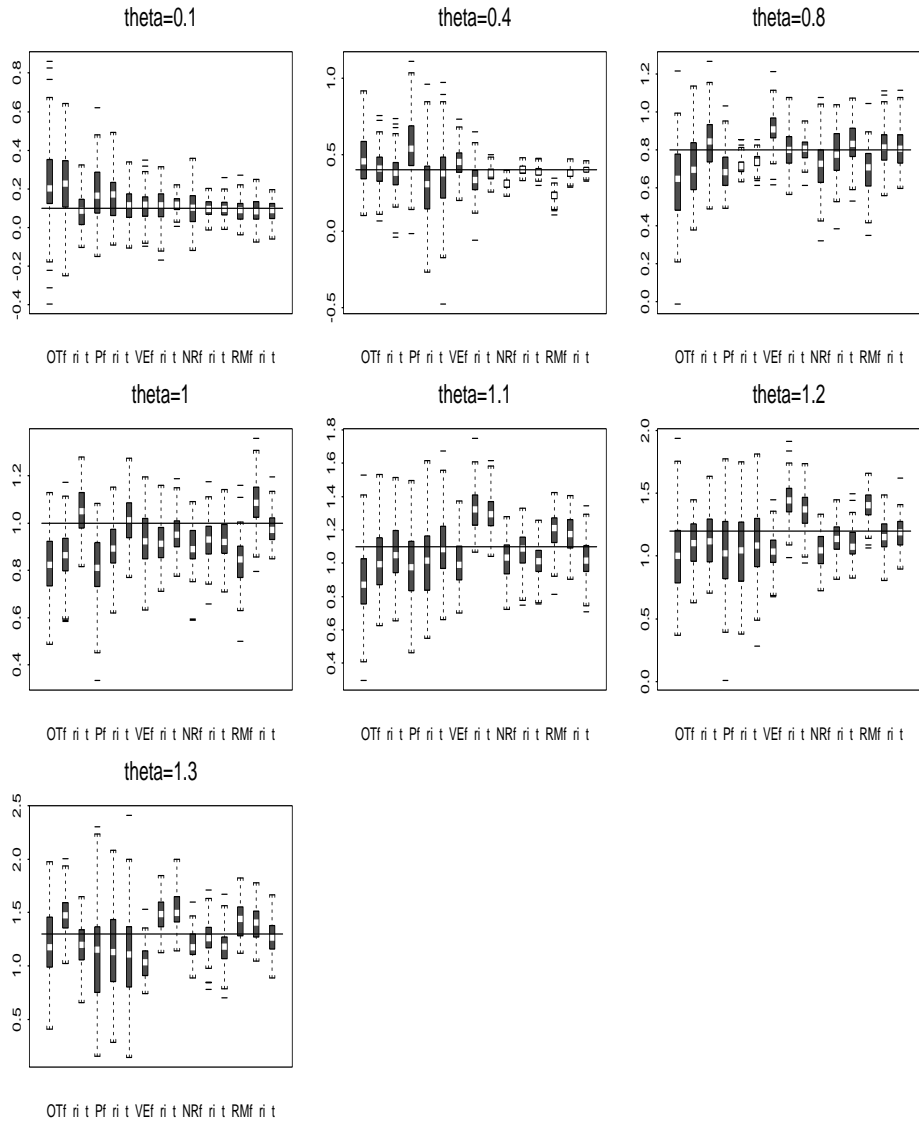
Figure 2b. Box-plots of simulated parameter estimates. See legend of Figure 2a. r=0.15

Figure 2c. Box-plots of simulated parameter estimates. See legend of Figure 2a. r=0.20

Table 1. Sample means and standard errors of parameter estimates when the interaction radius is $r = 0,10$. Each entry is based on 100 replicate simulations of n=50 events on the unit square. Lower case letters indicate the boundary condition: free, Ri=Ripley and toro=toroidal.

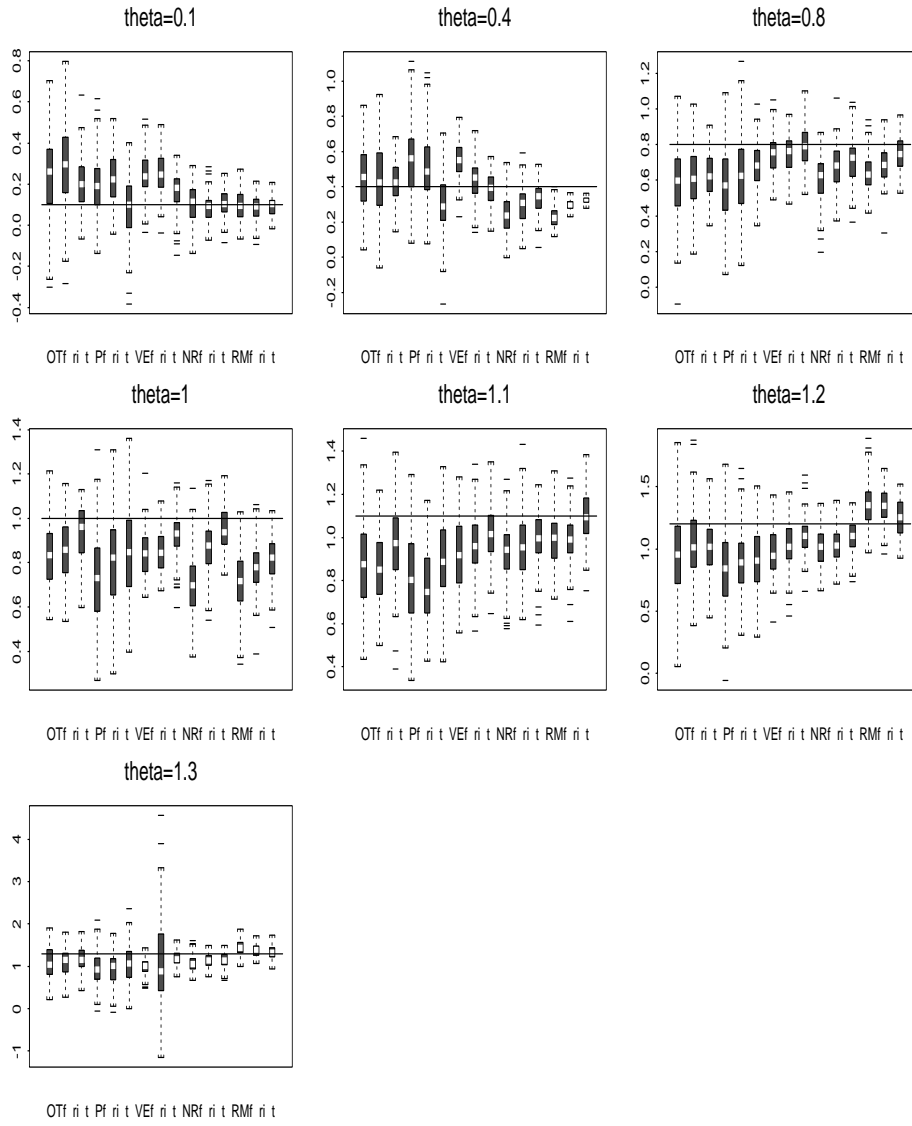| r=0.10	O-T method			Pent. method			V-E method			N-R method			R-M method		
	free	Ri	toro	free	Ri	toro	free	Ri	toro	free	Ri	toro	free	Ri	toro
θ	**Sample Means**														
0.1	0.247	0.231	0.091	0.193	0.176	0.087	0.112	0.125	0.122	0.095	0.097	0.098	0.093	0.091	0.096
0.4	0.434	0.427	0.407	0.541	0.308	0.345	0.456	0.319	0.376	0.305	0.403	0.393	0.231	0.378	0.404
0.8	0.645	0.714	0.835	0.703	0.715	0.732	0.913	0.809	0.805	0.704	0.791	0.807	0.710	0.803	0.824
1.0	0.831	0.873	1.041	0.810	0.912	1.013	0.919	0.920	0.946	0.910	0.935	0.931	0.847	1.079	0.979
1.1	0.914	0.979	1.093	0.973	1.007	1.081	1.013	1.315	1.291	1.004	1.073	0.993	1.215	1.183	1.032
1.2	0.997	1.101	1.103	1.031	1.046	1.097	1.035	1.416	1.335	1.053	1.143	1.103	1.392	1.194	1.197
1.3	1.124	1.445	1.213	1.093	1.106	1.148	1.056	1.496	1.531	1.197	1.292	1.197	1.431	1.393	1.245
θ	**Standard Errors**														
0.1	0.215	0.210	0.091	0.141	0.115	0.091	0.090	0.091	0.037	0.093	0.051	0.047	0.061	0.057	0.053
0.4	0.171	0.135	0.131	0.205	0.217	0.215	0.099	0.101	0.048	0.039	0.032	0.031	0.043	0.038	0.029
0.8	0.176	0.156	0.132	0.115	0.043	0.039	0.105	0.101	0.066	0.127	0.125	0.112	0.128	0.118	0.105
1.0	0.125	0.112	0.111	0.127	0.115	0.113	0.107	0.096	0.080	0.098	0.083	0.081	0.103	0.097	0.065
1.1	0.215	0.203	0.193	0.235	0.215	0.195	0.135	0.137	0.122	0.121	0.115	0.107	0.120	0.119	0.117
1.2	0.323	0.213	0.211	0.341	0.312	0.247	0.156	0.165	0.146	0.143	0.135	0.127	0.142	0.138	0.129
1.3	0.351	0.225	0.212	0.451	0.410	0.393	0.170	0.171	0.157	0.161	0.149	0.143	0.160	0.153	0.141

Table 2. Sample means and standard errors of parameter estimates when the interaction radius is $r = 0,15$.

r=0.15	O-T method			Pent. method			V-E method			N-R method			R-M method		
	free	Ri	toro	free	Ri	toro	free	Ri	toro	free	Ri	toro	free	Ri	toro

θ	Sample Means														
0.1	0.197	0.215	0.142	0.198	0.186	0.082	0.197	0.183	0.172	0.094	0.096	0.096	0.092	0.092	0.093
0.4	0.431	0.416	0.412	0.515	0.412	0.319	0.511	0.392	0.431	0.304	0.393	0.409	0.214	0.341	0.431
0.8	0.613	0.705	0.841	0.609	0.674	0.705	0.819	0.812	0.811	0.703	0.793	0.805	0.695	0.849	0.907
1.0	0.813	0.912	1.010	0.805	0.845	0.906	0.905	0.896	0.915	0.921	0.935	0.963	0.896	0.945	1.031
1.1	0.887	0.946	1.035	0.874	0.885	0.948	1.193	1.203	1.195	0.973	0.987	0.995	0.994	0.998	1.051
1.2	0.944	0.997	1.102	1.005	1.045	1.103	1.298	1.305	1.399	1.047	1.031	1.034	1.314	1.293	1.227
1.3	1.034	1.125	1.204	1.091	1.112	1.131	1.423	1.397	1.430	1.092	1.141	1.195	1.443	1.348	1.338

θ	Standard Errors														
0.1	0.212	0.205	0.146	0.150	0.141	0.115	0.102	0.052	0.036	0.091	0.050	0.045	0.061	0.060	0.055
0.4	0.202	0.165	0.108	0.210	0.231	0.212	0.103	0.096	0.047	0.039	0.037	0.031	0.048	0.035	0.028
0.8	0.176	0.135	0.126	0.212	0.195	0.118	0.105	0.037	0.071	0.131	0.129	0.125	0.127	0.121	0.113
1.0	0.146	0.131	0.127	0.210	0.196	0.121	0.104	0.096	0.091	0.103	0.095	0.091	0.113	0.093	0.091
1.1	0.215	0.201	0.153	0.235	0.221	0.198	0.135	0.131	0.118	0.125	0.121	0.119	0.125	0.122	0.119
1.2	0.345	0.303	0.246	0.319	0.312	0.251	0.158	0.162	0.135	0.148	0.137	0.132	0.157	0.138	0.118
1.3	0.431	0.397	0.353	0.425	0.418	0.401	0.177	0.182	0.152	0.159	0.152	0.151	0.191	0.153	0.139

Table 3. Sample means and standard errors of parameter estimates when the interaction radius is $r = 0,20$.

r=0.20	O-T method			Pent. method			V-E method			N-R method			R-M method		
	free	Ri	toro	free	Ri	toro	free	Ri	toro	free	Ri	toro	free	Ri	toro

θ	Sample Means														
0.1	0.240	0.273	0.184	0.210	0.221	0.082	0.253	0.231	0.165	0.092	0.093	0.093	0.090	0.091	0.093
0.4	0.450	0.443	0.415	0.539	0.495	0.324	0.551	0.453	0.397	0.241	0.295	0.335	0.221	0.298	0.321
0.8	0.595	0.593	0.625	0.605	0.614	0.693	0.719	0.771	0.789	0.615	0.693	0.705	0.630	0.710	0.747
1.0	0.841	0.855	0.931	0.741	0.793	0.845	0.810	0.839	0.921	0.710	0.845	0.947	0.708	0.793	0.810
1.1	0.872	0.879	0.947	0.793	0.815	0.897	0.915	0.986	1.023	0.941	0.943	0.998	0.983	0.995	1.108
1.2	0.936	0.979	0.998	0.847	0.895	0.913	1.005	1.039	1.103	1.009	1.015	1.074	1.334	1.321	1.253
1.3	1.013	1.115	1.197	0.915	0.945	1.041	1.027	1.093	1.154	1.051	1.123	1.147	1.451	1.382	1.324

θ	Standard Errors														
0.1	0.230	0.211	0.135	0.151	0.134	0.127	0.105	0.104	0.098	0.091	0.085	0.073	0.065	0.058	0.047
0.4	0.215	0.197	0.109	0.198	0.197	0.201	0.107	0.107	0.098	0.102	0.098	0.085	0.058	0.032	0.019
0.8	0.178	0.167	0.132	0.201	0.195	0.119	0.103	0.102	0.101	0.131	0.130	0.121	0.115	0.103	0.102
1.0	0.153	0.136	0.131	0.203	0.201	0.202	0.104	0.095	0.092	0.134	0.128	0.112	0.121	0.098	0.095
1.1	0.198	0.185	0.176	0.218	0.212	0.205	0.142	0.139	0.126	0.141	0.139	0.129	0.128	0.125	0.123
1.2	0.351	0.298	0.255	0.325	0.301	0.247	0.181	0.173	0.148	0.158	0.138	0.132	0.199	0.143	0.128
1.3	0.398	0.299	0.301	0.441	0.412	0.395	0.183	0.179	0.165	0.171	0.163	0.159	0.183	0.161	0.146

3.5. Results and discussion

Tables 1, 2 and 3 give the results of the simulation study, expressed in terms of the *sample means*, $\overline{\theta}$ and *standard errors*, s_{θ}. These two statistics characterize the sampling distribution of the parameter estimates as noted in section 3.2 above as neither the theoretical nor asymptotic approximations of the parameter distribution are not known. The bias and efficiency of the estimation can only be assessed by means of a Monte Carlo approach. However, Bayesian procedures could also be used to approach the theoretical parameter distribution as in Mateu and Montes (1995).

The table values indicate that both stochastic approximation methods, Newton-Raphson (N-R) and Robbins-Monro (R-M), exhibited better results, in terms of bias and standard errors, than Ogata-Tanemura (O-T) and Penttinen (P) methods for cases of strong regularity ($\theta \leq 0,4$) and clustering ($\theta \geq 1,1$). The approximate maximum likelihood method based on virial expansions (V-E) exhibited substantial bias, particularly when θ is large; however, this is qualitatively predictable on theoretical grounds, since the adequacy of the approximation to the likelihood deteriorates as θ increases. Implementation of this approximation for any potential is straightforward if only low-order virial coefficients are required. This method is not suited for estimation in cases of strong interaction.

The O-T and P approximate maximum likelihood methods provided substantial negative bias for medium-to-large values of θ giving relatively large standard deviations for these values. These two approximations are based on the sparseness assumption and are not reliable methods for clustered patterns for which higher-order interactions become important.

Inspecting the standard errors in conjunction with the range of θ, we observe that approximate maximum likelihood using O-T and P methods, provide large standard deviations for small ($\theta \leq 0,4$) and large ($\theta \geq 1,1$) values of parameter θ and, in any case, they are much larger than those obtained with the other three methods.

For different values of θ, the choice of boundary condition becomes important. Generally, for any method and parameter values, the periodic boundary condition produced better results than Ripley's, and in turn they are better than those obtained with the free boundary condition. The N-R and R-M approximate maximum likelihood methods provided unbiased and efficient estimates for all ranges of parameter values, under periodic and Ripley's boundary condition. However, they provided biased estimates under the free boundary condition.

Comparing the behaviour of the bias and the standard errors of estimates among the three interaction radii, we observe that, under the same parameter value, method of estimation and boundary condition, biases and standard errors increased with r providing worse estimates for $r = 0,2$ compared with $r = 0,1$. For example, for $r = 0,2$ and using the R-M procedure with Ripley's correction, we get significant bias compared with the

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unbiased and efficient estimates obtained under the same conditions but with $r = 0,1$. Apart from this, all properties analysed above are also true for different interaction radii.

The Strauss process with $\theta > 1$ is not a good model for applications. It may result in simulation difficulties such as sensitivity to edge-conditions, poor mixing, etc (Gates and Westcott, 1986). Moreover, the spatial birth-and-death approach might not be the optimal choice. For the well-defined case $\theta < 1$, exact simulation of the Strauss process is possible using the Propp-Wilson algorithm (Moller, 1998; Kendall and Moller, 1999). Concerning edge-corrections, another possibility is to apply conditional simulation: simulate a point pattern into some guard area using the model conditional on the observed point pattern and then apply the guard area events in estimation.

4. CONCLUSIONS

The conclusions, taking into account the results of our simulation study, are the following:

1. Stochastic approximations generally provide better results, particularly for medium-to-large parameter values, than those based on the sparseness assumption which are not suited for estimation in cases of strong interaction. For small parameter values and small interaction radius ($r = 0,1$), the Ogata-Tanemura approximation exhibits very good results.

2. For small interaction radius and using stochastic approximations, Ripley's and periodic boundary condition provide unbiased and efficient estimates. This is not true when r increases.

3. Generally, periodic and Ripley's boundary condition exhibit better results than free boundary condition.

4. When r increases the biases and the standard errors increase for any method, particularly for the approximate maximum likelihood methods.

5. Finally, in cases of clustered processes we recommend to use stochastic approximations with Ripley's or toroidal boundary condition. In cases of strong regularity, we could also use approximations based on the sparseness assumption.

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